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Modeling "collisions"

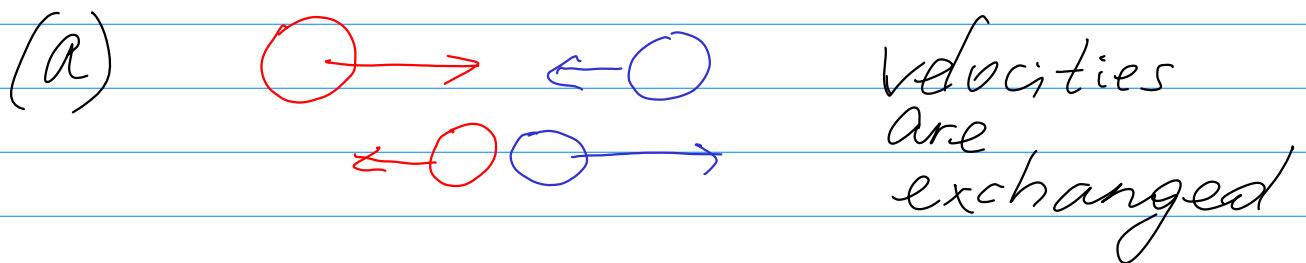
Rarefied gas, large number of particles.

$f(t, x, \xi)$ - density of particles at time t at point x with speed ξ .

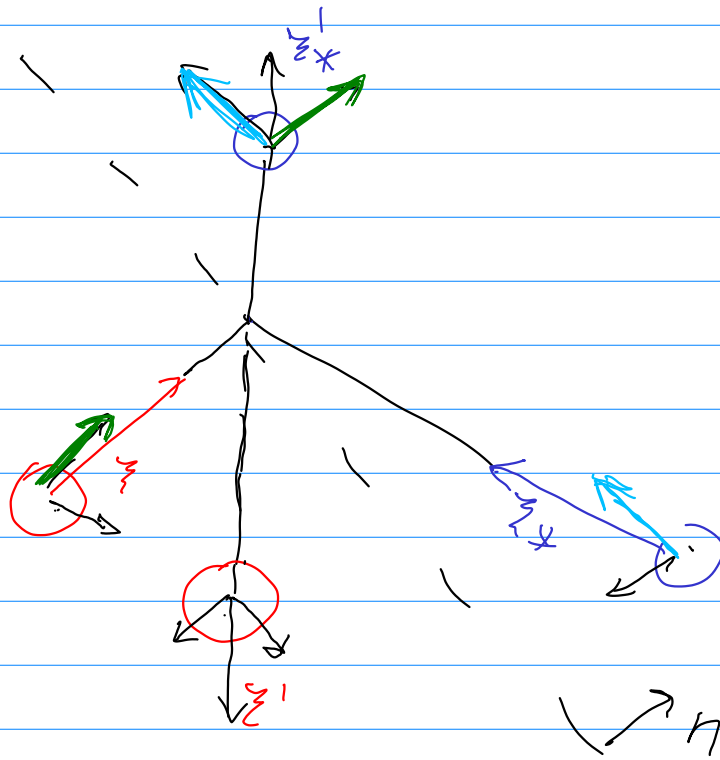
No collisions: $\partial_t f + \xi \cdot \nabla_x f = 0$

With collisions: $\partial_t f + \xi \cdot \nabla_x f = Q(f)$
↑
collision term

Elastic case: hard spheres
same radius



(b) in \mathbb{R}^3 - angle of collision is a unit vector n parallel to axis joining centers at collision



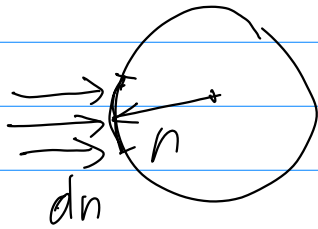
Components of velocities $\parallel n$ will be preserved, components $\perp n$ will be exchanged:

$$\begin{cases} \vec{z}' = \vec{z} - (n \cdot (\vec{z} - \vec{z}_*))n \\ \vec{z}'_* = \vec{z}_* + (n \cdot (\vec{z} - \vec{z}_*))n \end{cases} \quad \begin{array}{l} \text{Exchange} \\ \text{in } \perp n \text{ direction} \end{array}$$

new old velocities

$$\begin{cases} \vec{z}' \cdot n = \vec{z}_* \cdot n \\ \vec{z}'_* \cdot n = \vec{z} \cdot n \end{cases}$$

Deriving Boltzmann eqn:



Rate of collisions at an angle n is

$$|n \cdot v| dn$$

$\Rightarrow \dots \Rightarrow$

$$\xi - \xi_x$$

Loss term: $Q_-(f)(\xi) = \alpha \int_{\mathbb{R}^3} \int_{S^2} |n \cdot (\xi - \xi_x)| f(\xi) f(\xi_x) dn d\xi_x$

Elastic collision is reversible so

$$|n \cdot (\xi' - \xi'_x)| f(\xi') f(\xi'_x) = |n \cdot (\xi - \xi_x)| f(\xi) f(\xi_x)$$

\Rightarrow Gain term:

$$Q_+ f(\xi) = \alpha \int_{\mathbb{R}^3} \int_{S^2} |n \cdot (\xi - \xi_x)| f(\xi') f(\xi'_x) dn d\xi_x$$

Boltzmann eqn:

$$\partial_t f + \xi \cdot \nabla_x f = \int_{\mathbb{R}^3} \int_{S^2} (f' f'_x - f f_x) |n \cdot (\xi - \xi_x)| dn d\xi_x$$

H-theorem :

Rewrite Boltzmann as

$$\partial_t \mathcal{H}(t, x) + \nabla_x \cdot \mathcal{I}(t, x) = W(t, x)$$

where $\mathcal{H}(t, x) = \int_{\mathbb{R}^3} f \ln f d\xi = -S$

$$\mathcal{I} = \int_{\mathbb{R}^3} \xi f \ln f d\xi$$

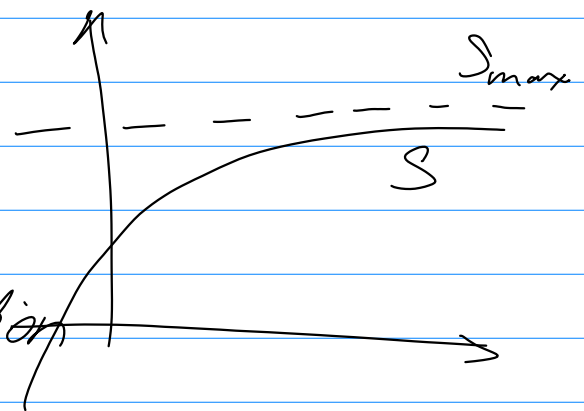
↑
Shannon-Boltzmann
entropy

$$W = \int_{\mathbb{R}^3} Q(f) \ln f d\xi$$

→ Since $W \leq 0$ and f decays at $|x| \rightarrow \infty$

$$\Rightarrow \boxed{\frac{d}{dt} \mathcal{H}(t) \leq 0}$$

Max entropy solution
is a Gaussian distribution



Beyond elastic case:

Granular flow

Smoluchowski
coagulation

$$v_1^* - v_2^* = -r(u_1 - u_2)$$



new vel.

old vel.

$$r = \begin{cases} 0 & \text{inelastic case} \\ 1 & \text{elastic case} \end{cases}$$

$$(u_1, u_2) \rightarrow (pu_1 + qu_2, qu_1 + pu_2)$$

$$p + q = 1, \quad r = 1 - 2p$$

$$K(v_1, v_2) = |v_1 - v_2|^n$$

collision rate

$$\lambda = 0 \Rightarrow f(v) \rightarrow \delta(v)$$

$$T = T_0 e^{-\lambda_2 t}$$

exp. decay of temperature

$$\lambda_n = 1 - p^n - q^n$$

$$V_1 + V_2 \rightarrow V^*$$

$$\text{mass } m^* = m_1 + m_2$$

$$\rightarrow p^* = p_1 + p_2$$

momentum

Kinetic energy
is dissipated

Elastic	Granular	Smoluchowski
Mass = const	Mass = const	Mean mass grows
$E_{\text{kinetic}} = \text{const}$	$E_{\text{kin}} \searrow$	$E_{\text{kin}} \searrow$ (cooling effect)
Reversible	Irreversible collision	